Tunable Proximity Effects and Topological Superconductivity in Semiconductor-Superconductor-Ferromagnetic Hybrid Nanowires

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Introduction

Arxiv preprint: Microscopic analysis of topological superconductivity in ferromagnetic hybrid nanowires arXiv:2011.06566



Alfredo Levy Yeyati *UAM*

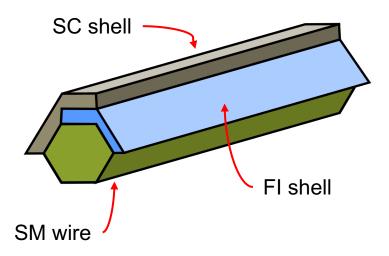


Elsa Prada



Yuval Oreg Weizmann Institute

Hybrid nanowire



Applications in:

- Spintronics
- Topological superconductivity

Motivation Outline

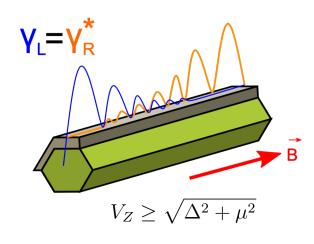
- MBS are topological subgap modes that can emerge at the ends of a superconductor/semiconductor nanowire

$$H = \left(\frac{\hbar^2 k^2}{2m^*} - \mu + V_Z \sigma_x + \vec{\alpha} \cdot (\vec{\sigma} \times \vec{k})\right) \tau_z - i\Delta\sigma_y \tau_y$$

Ingredients: Kinetic energy

Electrochemical potential Induced SO interaction superconductivity

Zeeman field:
$$V_Z=rac{1}{2}\mu_B g B_x$$



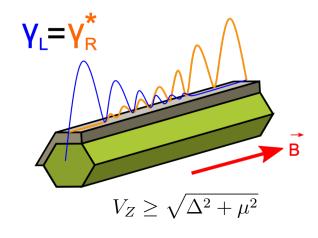
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Ingredients: Kinetic Electroenergy chemical potential SO interaction superconductivity

Zeeman field:
$$V_Z=rac{1}{2}\mu_B g B_x$$



- MBS may be useful for (topological) quantum computing.

Motivation Outline

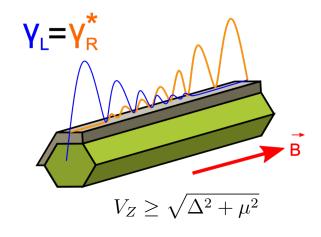
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- MBS may be useful for (topological) quantum computing.

But the magnetic field weakens the superconductivity and complicates the scaling of a QC...

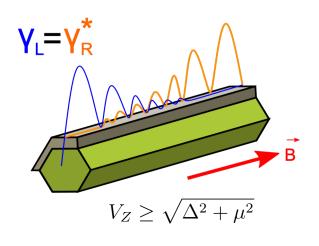
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- MBS may be useful for (topological) quantum computing.

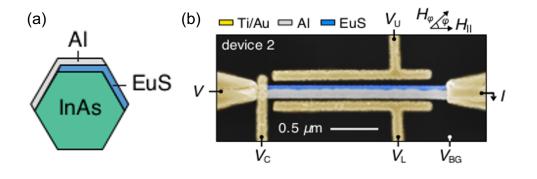
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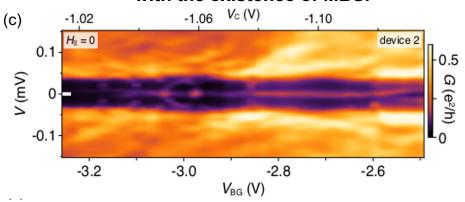
Is it possible to create MBS without a magnetic field?

Introduction

There is no need of an external magnetic field if it can be intrinsically incorporated. Recent experimental works show that it is possible to induce an exchange field in the nanowire by proximitizing an EuS layer to the heterostructure.

This device shows ZBP compatible with the existence of MBS.



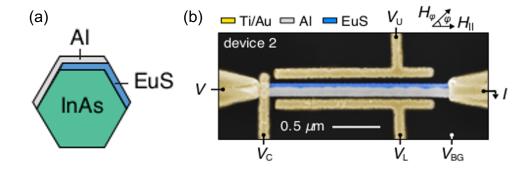


Refs. Y. Liu *et al.*, ACS App. Mat. **12**, 8780 (2020) Y. Liu *et al.*, Nano Lett. **20**, 456 (2020) S. Vaitiekėnas *et al.* Nat. Phys. 17, 43 (2021)

Introduction

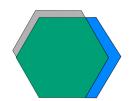
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(c) $\frac{-1.02}{0.1}$ $\frac{-1.06}{|W_c(V)|}$ $\frac{-1.10}{0.5}$ $\frac{1.10}{|W_c(V)|}$ $\frac{1.10}{0.5}$ $\frac{1.10}{|W_c(V)|}$ $\frac{1.10}{0.5}$ $\frac{1.10}{|W_c(V)|}$ $\frac{1.10}{0.5}$ $\frac{1.10}{|W_c(V)|}$ $\frac{1.10}{0.5}$ $\frac{1.10}{|W_c(V)|}$ $\frac{1.10}{0.5}$ $\frac{1.10}{|W_c(V)|}$ $\frac{1.10}{0.5}$ $\frac{1.10}{|W_c(V)|}$ $\frac{1.10}{0.5}$

Strikingly, other geometries show little or no induced magnetization



Could be the ZBP Majorana Bound States? Why these devices do not show ZBP? How is induced the magnetization?

ΑI

Introduction

Useful to understand the induced magnetization

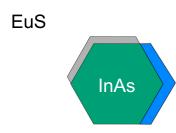
Useful to study the

phase diagram

- Realistic model
 - + Overlapping device
 - + Non-overlapping device
- Effective model
 - + Overlapping device
 - + Non-overlapping device
- Conclusions



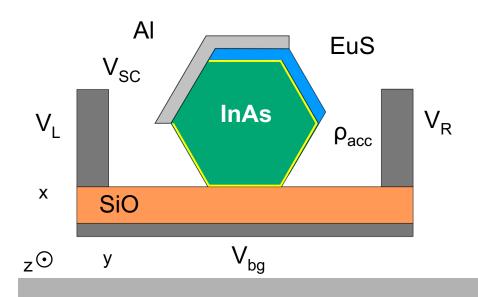
Overlapping device (shows ZBP)



Non-overlapping device (doesn't show ZBP)

We include in the Hamiltonian all the materials involved in the heterostructure using realistic parameters. We also include the self-consistent electrostatic environment.

$$H = \vec{k} \frac{\hbar^2}{2m_{\text{eff}}(\vec{r})} \vec{k} + E_{\text{F}}(\vec{r}) - e\phi(\vec{r}) + h_{\text{ex}}(\vec{r})\sigma_x + \Delta(\vec{r})\tau_x\sigma_x + \frac{1}{2} \left[\vec{\alpha}(\vec{r}) \cdot (\vec{\sigma} \times \vec{k}) + (\vec{\sigma} \times \vec{k}) \cdot \vec{\alpha}(\vec{r}) \right]$$



InAs		
m _{eff}	0.023m ₀	
E _F	0	
h _{ex}	0	
Δ	0	
$\alpha_{_{R}}$		

	EuS	
m _{eff}	0.3m ₀	
E _F	0.7eV	
h _{ex}	0.1eV	
Δ	0	
α_{R}	0	

Al		
m _{eff}	m ₀	
E _F	-10eV	
h _{ex}	0	
Δ	0.23meV	
α_{R}	0	

$$\vec{\alpha}(\vec{r}) = \vec{\alpha}_{\rm int} + \frac{eP_{\rm fit}^2}{3} \left[\frac{1}{\Delta_{\sigma}^2} - \frac{1}{(\Delta_{\rm g} + \Delta_{\rm soff})^2} \right] \vec{\nabla} \phi(\vec{r})$$

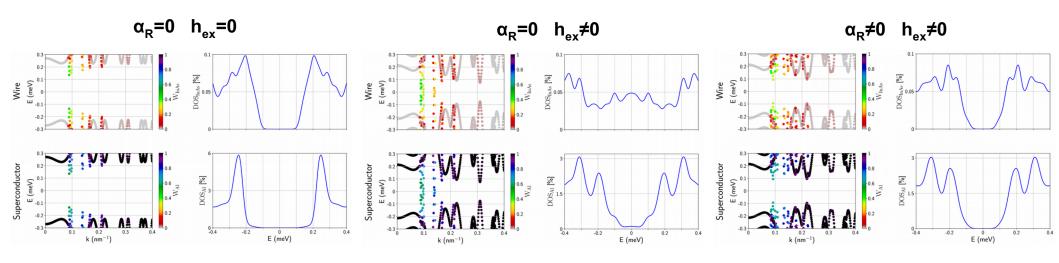
Model

Results

- Overlapping device
- Non-overlapping device

We compute the energy spectrum versus the momentum k_z for the **overlapping device** fixing all the gates to V_i =0. From there, we also compute the DOS. We perform three different simulations.

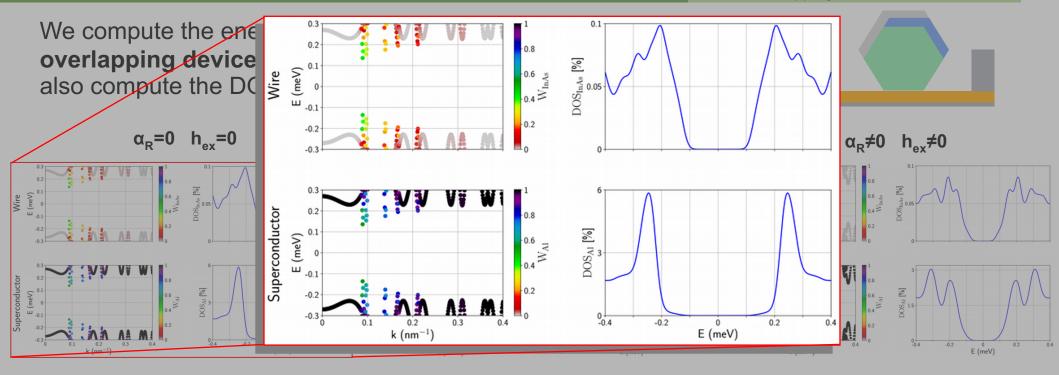




Model

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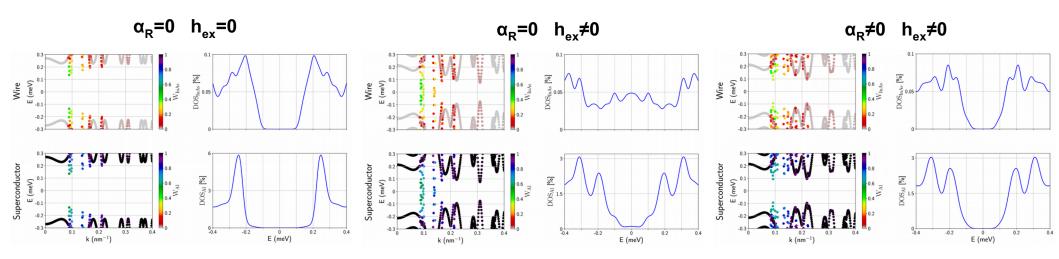


Model Results

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Induced gap of 0.2 meV

Model

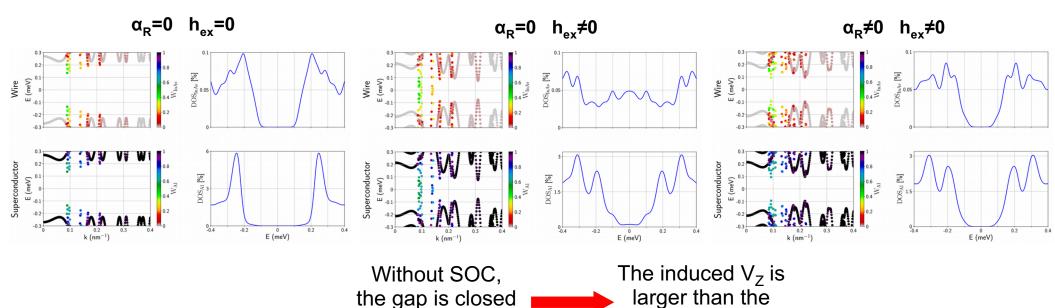
induced Δ

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in the NW



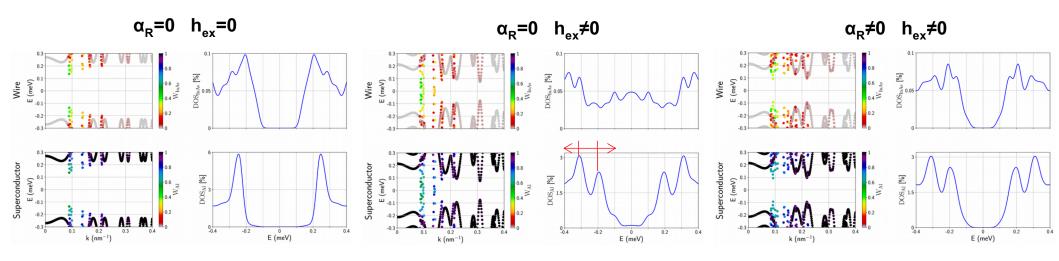


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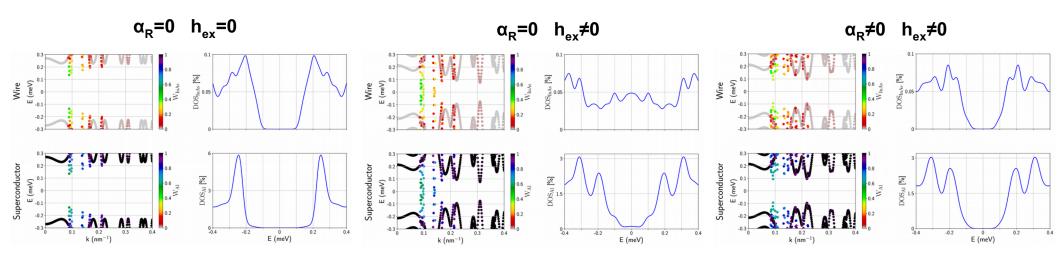
A small exchange field of 0.06meV is also induced in the SC, as previous experiments showed

Model Results

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- Non-overlapping device

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Signature of topological phase transition



With SOC, the gap reopens

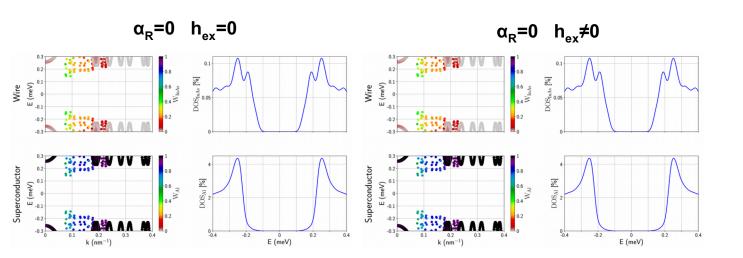
Model

Results

- Overlapping device
- Non-overlapping device

We perform exactly the same simulations but for the **non-overlapping device**.



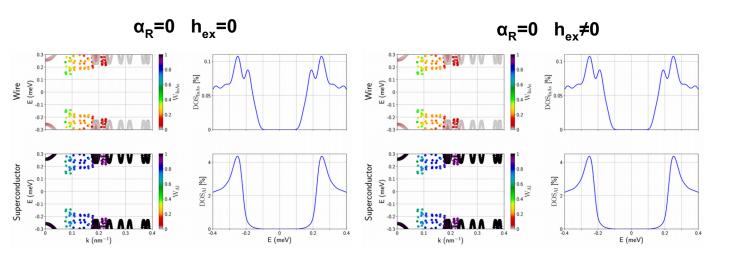


Model **Results**

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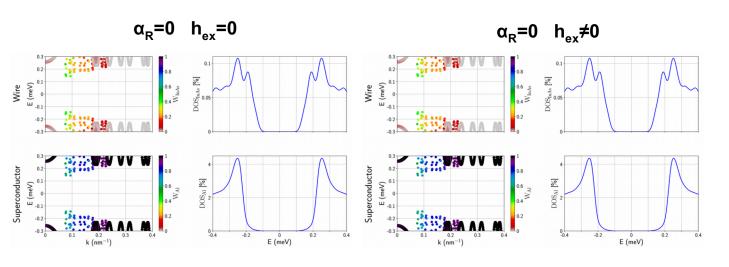
For the non-overlapping device, the induced exchange field seems not to be large enough to close the gap

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There is no topological phase in the non-overlapping device, at least for this gate voltage



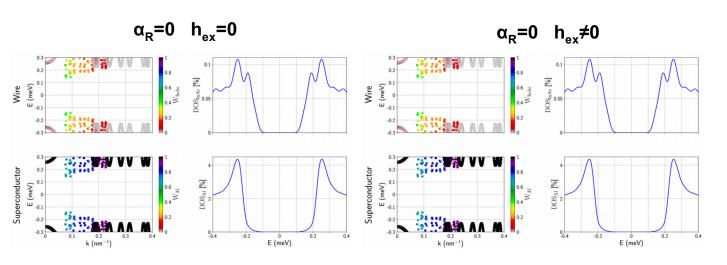
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Topological phase diagram?



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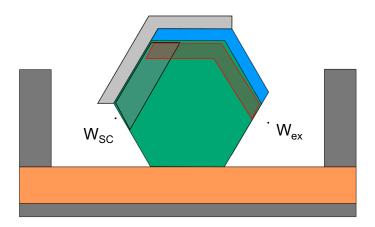
Model

Results

- Overlapping device
- Non-overlapping device

We "integrate out" the Al and the EuS, and we directly include the proximity effects into the InAs nanowire in an effective way. This reduces the computational cost and allows to find the phase diagram.

There is only a superconducting pairing of Δ =0.2meV in this proximitizing region (W_{SC}=30nm), as well as an exchange field of h_{ex} =0.06meV



There is an exchange field of h_{ex}^(EuS)=0.1eV in this proximitizing region (W_{ex}=1nm)

$$H = \frac{\hbar^2 k^2}{2m_{\text{eff}}} + E_{\text{F}} - e\phi(\vec{r}) + h_{\text{ex}}(\vec{r})\sigma_x + \Delta(\vec{r})\tau_x\sigma_x + \frac{1}{2} \left[\vec{\alpha}(\vec{r}) \cdot (\vec{\sigma} \times \vec{k}) + (\vec{\sigma} \times \vec{k}) \cdot \vec{\alpha}(\vec{r}) \right]$$

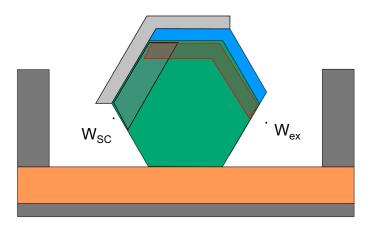
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There is an exchange field of h_{ex}^(EuS)=0.1eV in this proximitizing region (W_{ex}=1nm)

We compute the induced magnetization and superconductivity. We choose W_{SC} and W_{ex} in such a way to reproduce (roughly) the same behaviour as in the realistic model.

-2.5

100

25

100

 $h_{\rm ex}^{(EuS)}$ (meV)

 $h_{\rm ex}^{(EuS)}$ (meV)

Phase diagram vs V_{bg} (fixing V_L =0 and V_R =-4V) for an **overlapping** device with direct-induced magnetization

-2

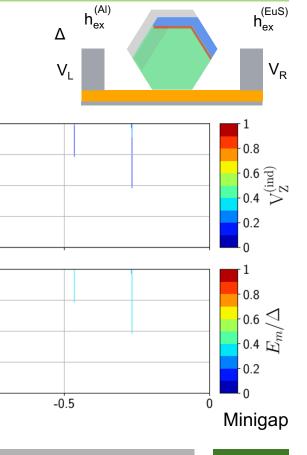
-1.5

 $V_{bq}\left(\mathsf{V}\right)$

Model

Results

- Overlapping device
- Non-overlapping device



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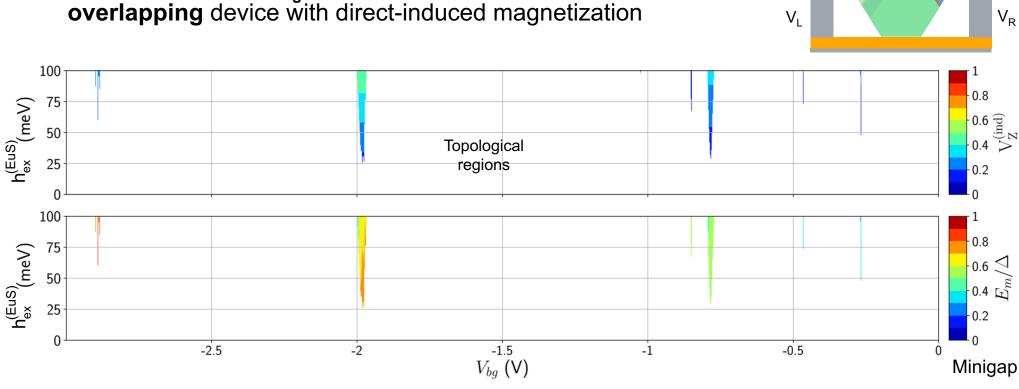
Model

Results

- Overlapping device

h_{ex}

- Non-overlapping device



h_{ex}(EuS)

-2.5

100

25

100

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Phase diagram vs V_{bg} (fixing V_L =0 and V_R =-4V) for an **overlapping** device with direct-induced magnetization

-2

Induced V₇~0.4meV

Induced gap of 0.6Δ

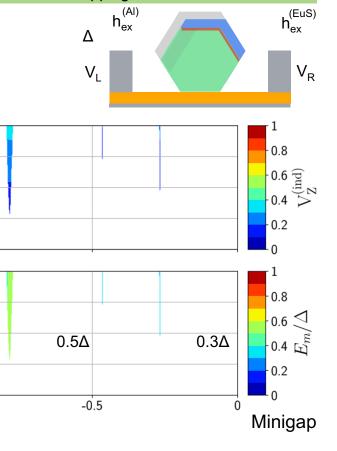
-1.5

 $V_{bq}\left(\mathsf{V}\right)$

Model

Results

- Overlapping device
- Non-overlapping device



100

100

 $h_{\rm ex}^{(EuS)}$ (meV)

 $h_{\rm ex}^{\rm (EuS)}$ (meV)

Phase diagram vs V_R (fixing $V_L=0$ and $V_{bg}=-2V$) for an **overlapping** device with direct-induced magnetization

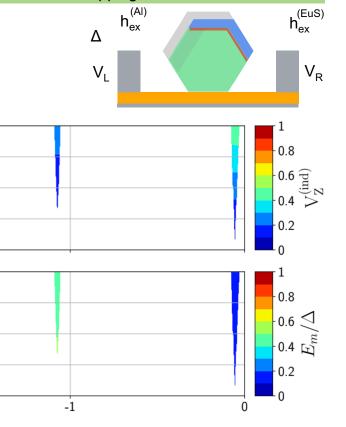
-3

 V_R (V)

Model Results

- Overlapping device

- Non-overlapping device



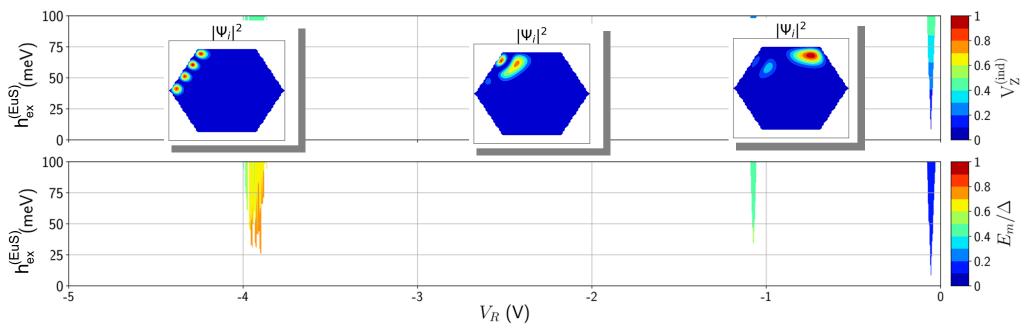
Phase diagram vs V_R (fixing $V_L=0$ and $V_{bg}=-2V$) for an **overlapping** device with direct-induced magnetization

Model

Results

- Overlapping device
- Non-overlapping device



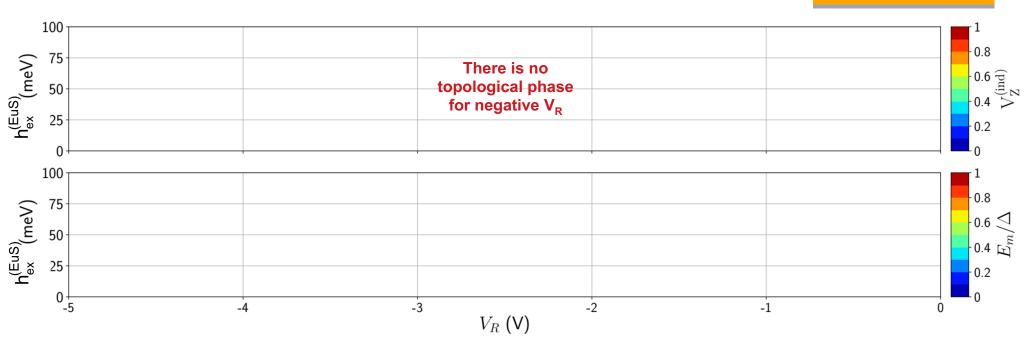


Model Results

- Overlapping device
- Non-overlapping device

Phase diagram vs V_R (fixing V_L =0 and V_{bg} =-2V) for a **non-overlapping** device with direct-induced magnetization





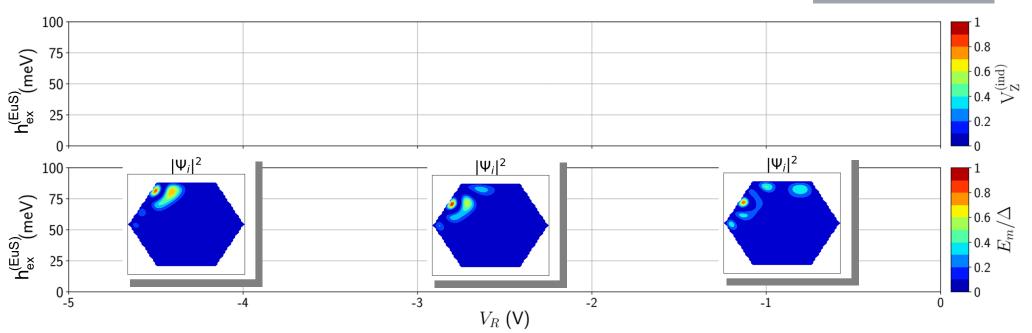
Model

Results

- Overlapping device
- Non-overlapping device

Phase diagram vs V_R (fixing V_L =0 and V_{bg} =-2V) for a **non-overlapping** device with direct-induced magnetization





Conclusions and outlook

Conclusions

- InAs/Al/EuS heterostructures intrinsically incorporates the effect of a Zeeman field large enough so that they can support MBS.
- Only some specific geometries give rise to MBS, because the wavefuntion needs to be close to the EuS-InAs and Al-InAs interfaces at the same time. The strength of the proximity effects can be controlled by the gates.

Reference

- Microscopic analysis of topological superconductivity in ferromagnetic hybrid nanowires, Samuel D. Escribano, Elsa Prada, Yuval Oreg and Alfredo Levy Yeyati, arXiv:2011.06566 (2020).

For any question or inquire, don't hesitate to contact me via email at **samuel.diaz@uam.es**, thank you for your attention!

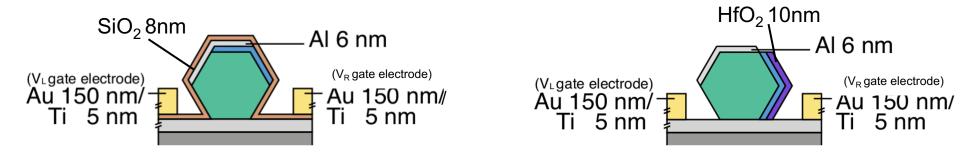
Supplementary Material

A: Effective Model

Induced superconductivity
Induced Zeeman field

The electrostatic potential is determined self-consistently (in the Thomas-Fermi approximation) using the Poisson equation. The electrostatic environment is taken into account through the dielectric permittivity.

$$\vec{\nabla}(\epsilon(\vec{r}) \cdot \vec{\nabla}\phi(\vec{r})) = \rho(\vec{r})$$



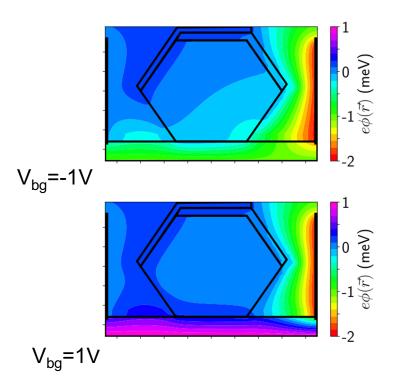
A recent experiment shows that there is an accumulation layer at the InAs-EuS interface similar to the one of the free facets. Thus, we include the same accumulation layer ρ_{acc} in the nanowire facets that are not in contact with Al. Additionally, we simulate the InAs-Al band bending imposing V_{SC} as boundary condition on the Al.

Model

Electrostatic potential

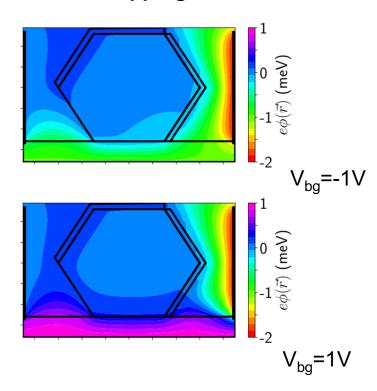
Induced superconductivity Induced Zeeman field

Overlapping device



As the back-gate voltage is increased, the wavefunction is pushed towards the bottom of the wire.

Non-overlapping device

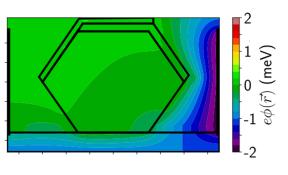


Model

Electrostatic potential

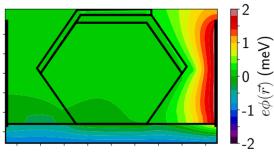
Induced superconductivity Induced Zeeman field

Overlapping device



$$V_R = -2V$$

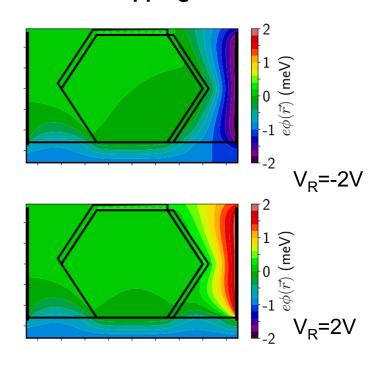
As the right-gate voltage is increased the wavefunction is pushed towards the EuS.



The proximity effects, both with Al and EuS can thus be controlled by the gates.

$$V_R=2V$$

Non-overlapping device



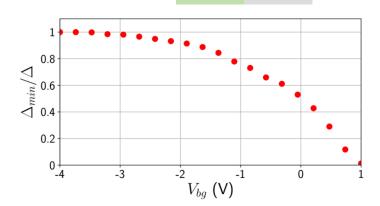
Model

To describe the superconductivity inside the semiconductor, one would need, in principle, to include the superconducting layer also at a tight-binding level.

The SC is described as a metallic region (with a band-offset of -10eV) with a paring amplitude Δ



One can obtain the spectra of the system for different gates, and from there, the DOS in the wire and the induced gap (the minimum gap Δ_{min}).



-0.2eV

10eV

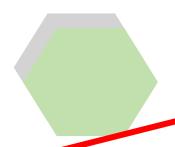
InAs

Αl

To describe the superconductivity inside the semiconductor, one would need, in principle,

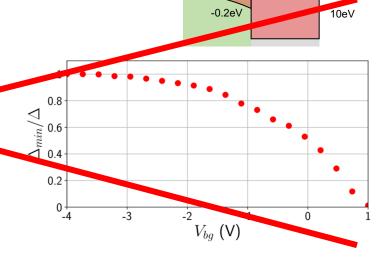
to include the superconducting layer also at a tight-binding level.

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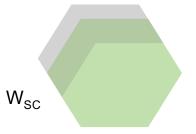




Unfortunately, this is not computationally affordable.

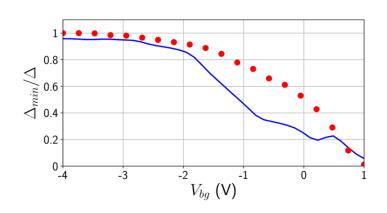
A different approach to include the proximity effect in the wire is to assume that a region of width W_{SC} close to the InAs/Al interface is characterized by a paring amplitude Δ .

The SC is described as a hard wall (not included in the TB)



 Δ is present only in this region

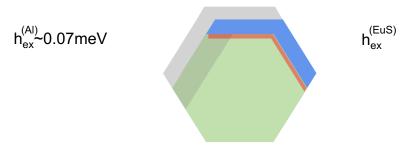
It is possible to do the same for this system (blue line).



Using W_{SC}=30nm we predict a similar behaviour.

It is not clear how the magnetization induced by the EuS influences the state of the nanowire. There are two possible scenarios, which could be complementary.

Model 1: direct-induced magnetization



The EuS **directly** induces an exchange field ($h_{ex}^{(EuS)}$) in the InAs. Because the EuS is an insulator, the proximitized region is small (1nm), but with a large exchange field. In addition, it is known that there is a small exchange field ($h_{ex}^{(Al)}$ =0.07meV) in the Al due to the Al/EuS interface.

Model 2: indirect-induced magnetization



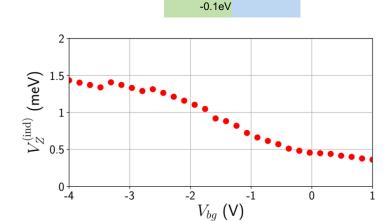
The EuS induces an exchange field $h_{ex}^{(Al)}$ in the InAs through the Al layer in an **indirect** way. The exchange field induced in the SC due to the Al-EuS interface is indeed, for whichever reason, larger than Δ . The spinorbit coupling opens a gap even if the Clogston limit is reached.

To show that the first model is also plausible, let us describe first the EuS at a tight-binding level as well.

The EuS is described as an insulating region (with a band-offset of 0.8eV) characterized by a large exchange field h_{ex} (with a Zeeman splitting of 0.1eV)



One can obtain the spectra of the system for different gates, and from there, the induced magnetization in the wire.

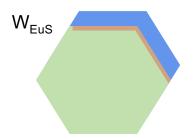


0.1eV

0.8eV

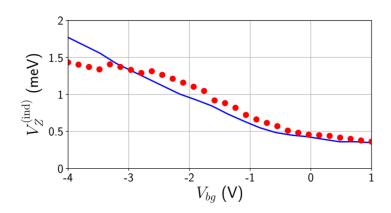
Although it is (computationally) affordable to include the EuS at a tight-binding layer, let us describe it as a proximitized region close to the InAs-EuS interface, as we did for the Al.

The EuS is described as a hard wall (not included in the TB)



h_{ex} is present only in this region

It is possible to do the same for this system (blue line).



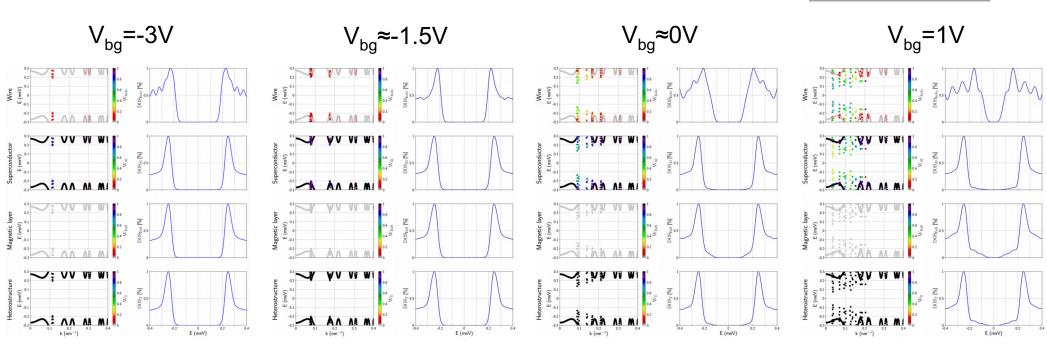
Using W_{EuS} =1nm and h_{ex} \simeq 100meV we predict a similar behaviour.

Supplementary Material

B: DOS vs V_{bg}

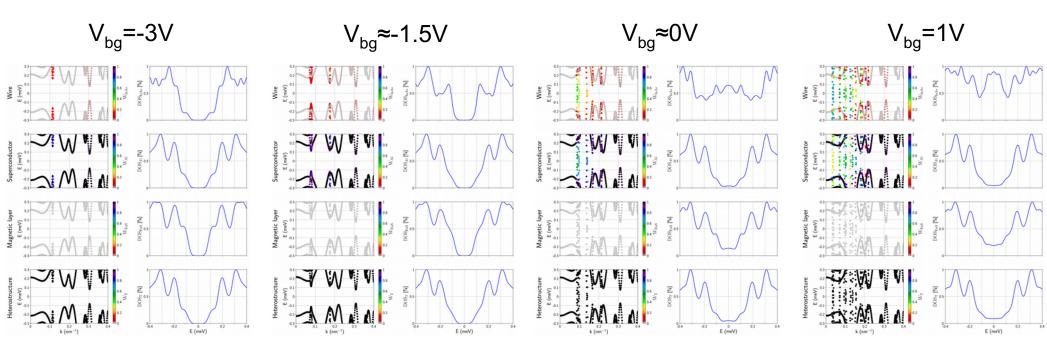
DOS vs V_{bg} for the **overlapping device** with h_{ex} =0 and α_{R} =0.





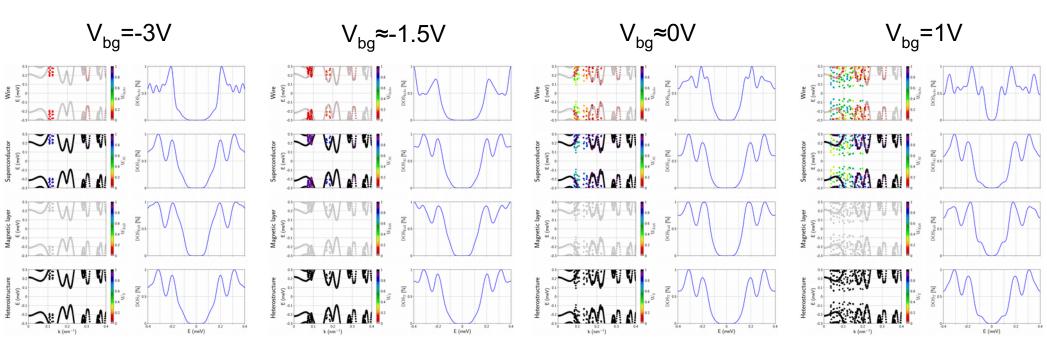
DOS vs V_{bq} for the **overlapping device** with $h_{ex} \neq 0$ and $\alpha_{R} = 0$.



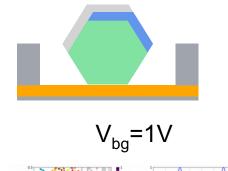


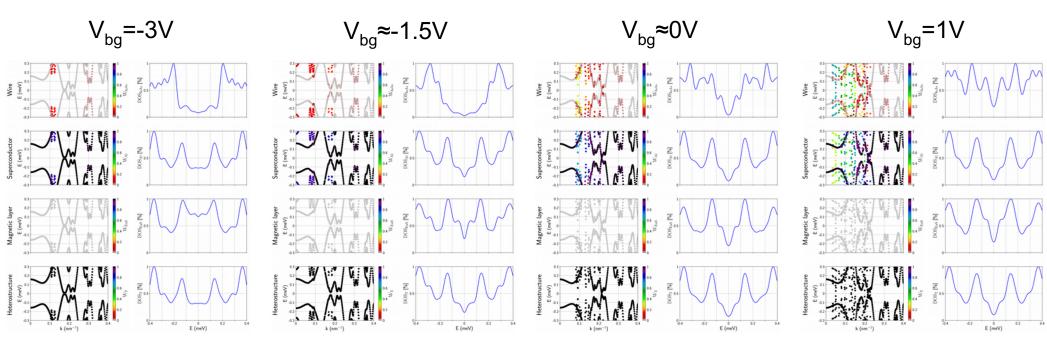
DOS vs V_{bg} for the **overlapping device** with $h_{ex} \neq 0$ and $\alpha_{R} \neq 0$.

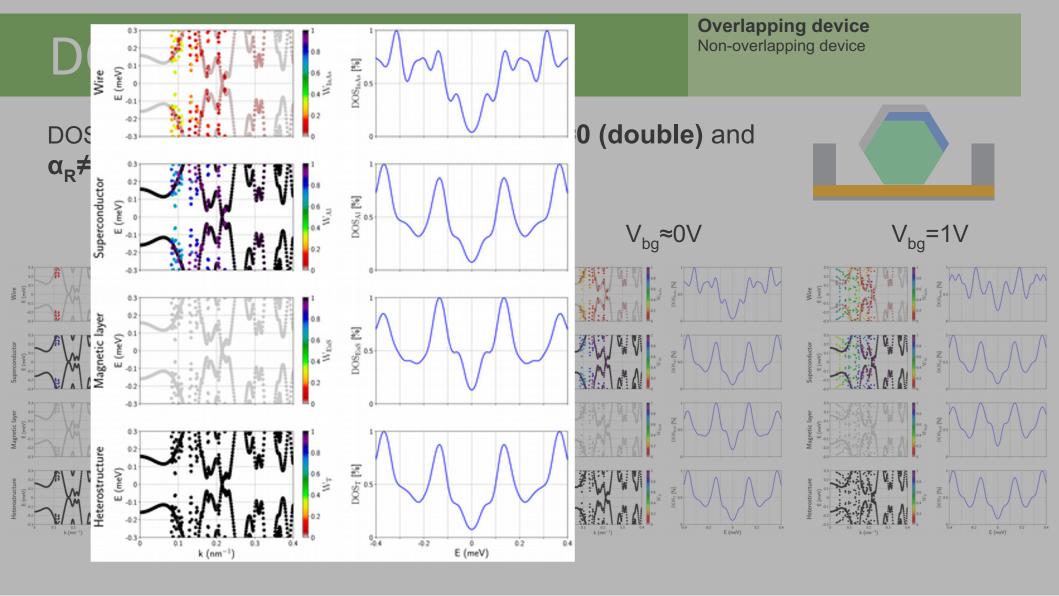




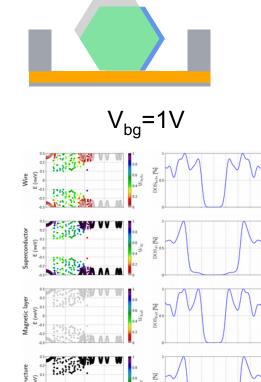
DOS vs V_{bg} for the **overlapping device** with $h_{ex}\neq 0$ (double) and $\alpha_R\neq 0$.

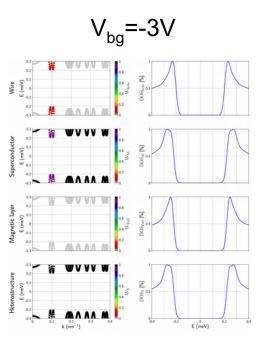


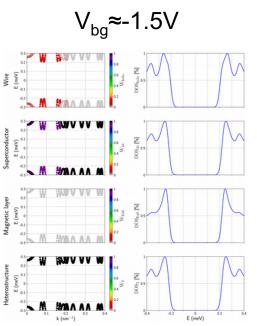


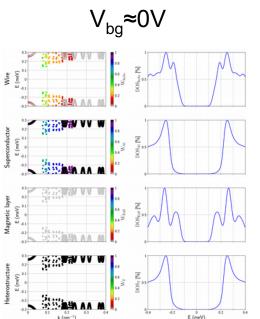


DOS vs V_{bq} for the **non-overlapping device**.







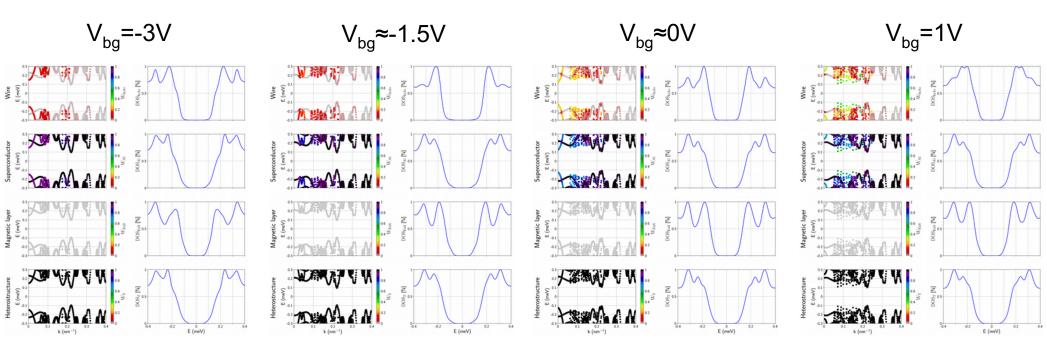


Supplementary Material

C: 4-facets geometry

DOS vs V_{bg} for the **4-facets device** with $h_{ex} \neq 0$ and $\alpha_R = 0$.





Supplementary Material

D: Extended phase diagram for the non-overlapping

Effective model

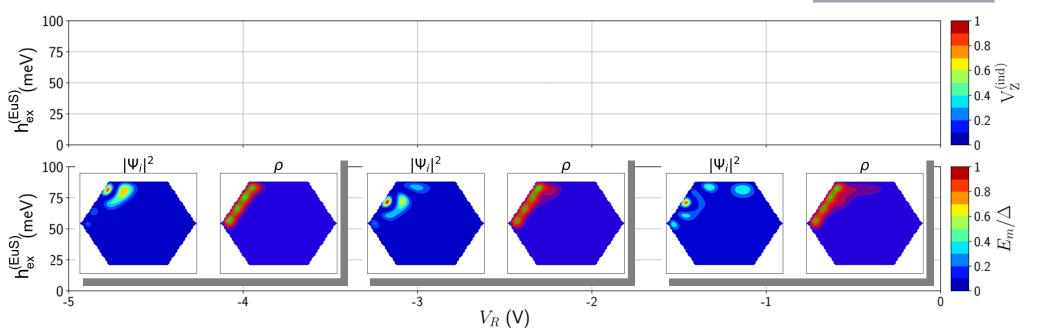
Model

- Results
 - Non-overlapping device

- Overlapping device

Phase diagram vs V_R (fixing $V_L=0$ and $V_{bg}=-2V$) for a **non-overlapping** device with direct-induced magnetization





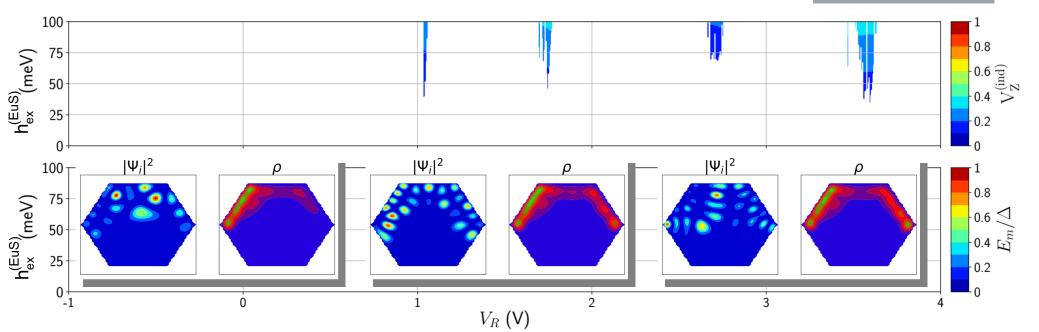
Effective model

Model

- Results
 - Overlapping deviceNon-overlapping device

Phase diagram vs V_R (fixing V_L =0 and V_{bg} =-2V) for a **non-overlapping** device with direct-induced magnetization





Effective model

Model **Results**

- Overlapping device
- Non-overlapping device

Phase diagram vs V_R (fixing V_L =0 and V_{bg} =-2V) for a **non-overlapping** device with direct-induced magnetization



